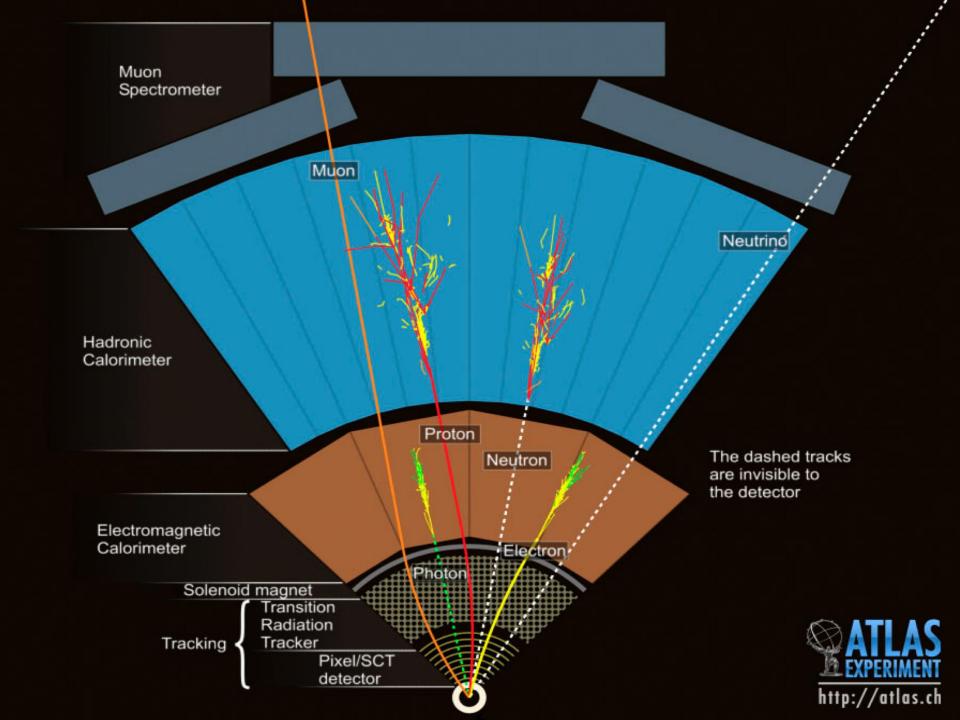
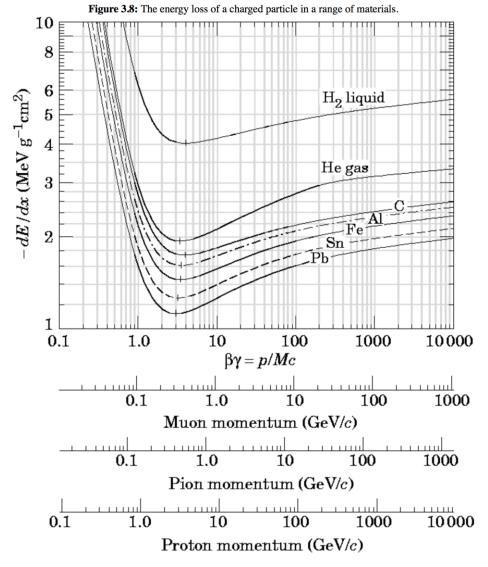
Measuring the Gluon Spin Distribution at Small-x (Part 2: Detector)

Mickey Chiu BROOKHAVEN NATIONAL LABORATORY



Ionization Loss of Charged Particles in

Matter



 Bethe-Bloch Equation... Coulomb Interaction with the electrons in the atoms

13.1 Photons in matter

(Overview)

Rayleigh scattering

- Coherent, elastic scattering of the entire atom (the blue sky)
- $-\gamma$ + atom $\rightarrow \gamma$ + atom
- dominant at λ_{γ} >size of atoms

Compton scattering

- Incoherent scattering of electron from atom
- $-\gamma + e_{\text{bound}} \rightarrow \gamma + e_{\text{free}}$
- possible at all $E_{\gamma} > min(E_{bind})$
- to properly call it Compton requires $E_{\gamma} >> E_{bind}(e^{-})$ to approximate free e^{-}

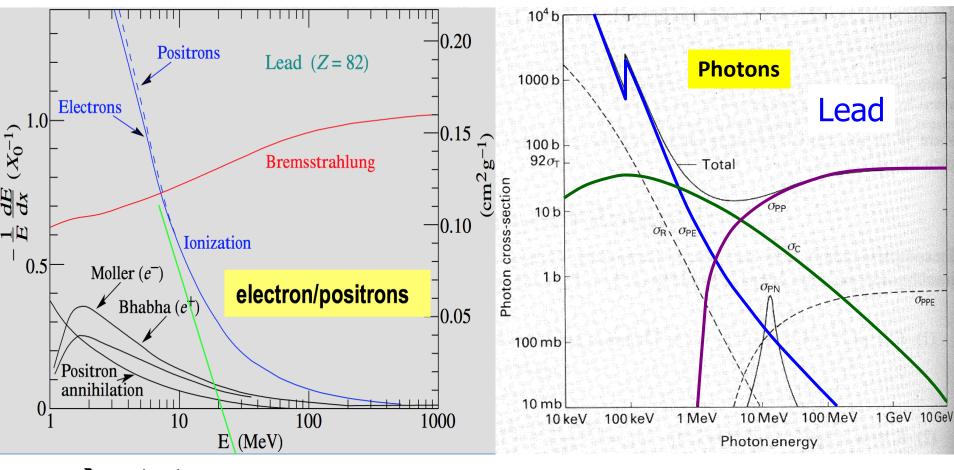
Photoelectric effect

- absorption of photon and ejection of single atomic electron
- $-\gamma + atom \rightarrow \gamma + e_{free}^- + ion$
- possible for E_{γ} < max(E_{bind}) + $\delta E(E_{atomic-recoil}$, line width) (just above k-edge)

Pair production

- absorption of γ in atom and emission of e^+e^- pair
- Two varieties:
 - γ + nucleus \rightarrow e⁺ + e⁻ + nucleus (more momentum transfer to nucleus \rightarrow dominates)
 - γ + Z atomic electrons \rightarrow e⁺ + e⁻ + Z atomic electrons
 - both summarised via: g + g(virtual) → e⁺ + e⁻
- Needs $E\gamma > 2m_e c^2$
- Nucleus has to recoils to conserve momentum → coupling to nucleus needed → strongly Z-dependent crossection

Electromagnetic Interactions in Matter



- R → Rayleigh
- PE → Photoeffect
- $C \rightarrow Compton$

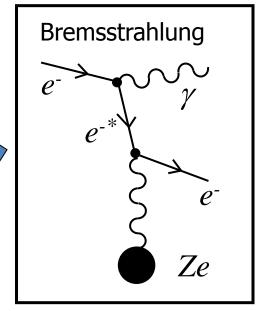
- PP → Pair Production
- PPE → Pair Production on atomic electrons
- PN → Giant Photo-Nuclear dipole resonance

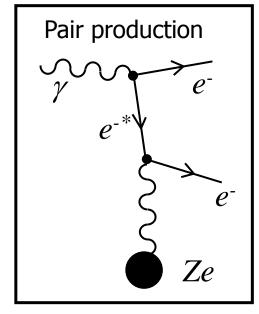
13.1 Photons in matter

(Note on Pair Production)

Compare pair production with Bremsstrahlung

Typical Lenth = Radiation Length X0





Typical Lenth = Pair Production Length L0



- Very similar Feynman Diagram
- Just two arms swapped

L0=9/7 X0

Electromagnetic showers

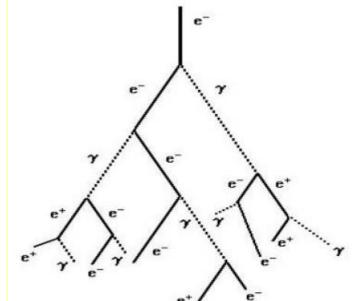
(A) Bremsstrahlung

radiation of real photons in the Coulomb field of nuclei

$$-\frac{dE}{dx} = 4\alpha \cdot N_A \cdot \frac{Z^2}{A} \cdot z^2 \cdot \left(\frac{1}{4\pi\varepsilon_0} \frac{e^2}{mc^2}\right) \cdot \ln\left(\frac{183}{Z^{\frac{1}{3}}}\right) \cdot E$$

(B) Pair production

needs additional mass for momentum conservation

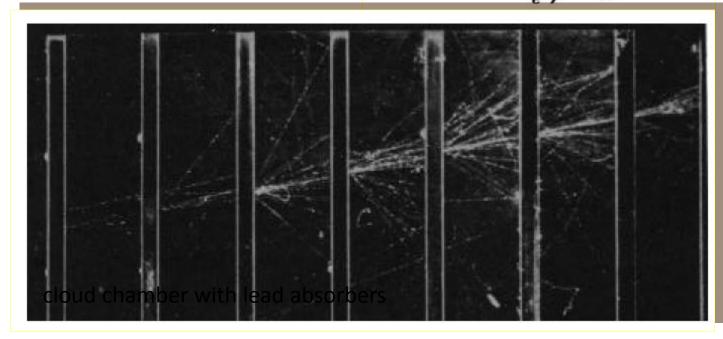


$$X_0 = \frac{A}{4\alpha \cdot N_A Z^2 r_e^2 \ln\left(\frac{183}{Z^{\frac{1}{3}}}\right)}$$

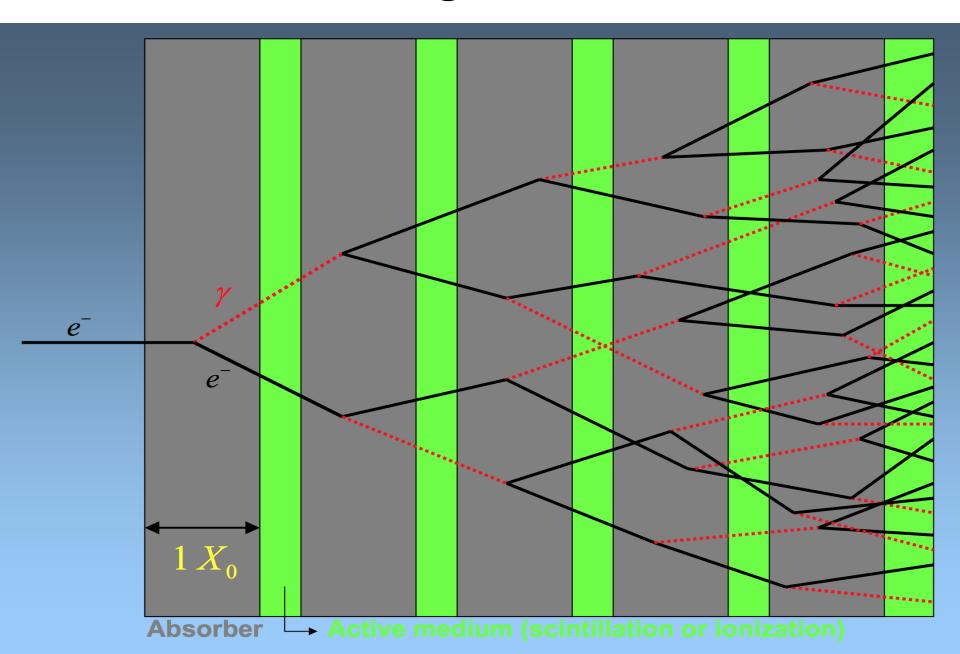
$$\lambda_{pair} = \frac{7}{9} X_0$$

Moliere radius

$$R_{M} = X_{0} \frac{21MeV}{E_{C}}$$



Electromagnetic Shower



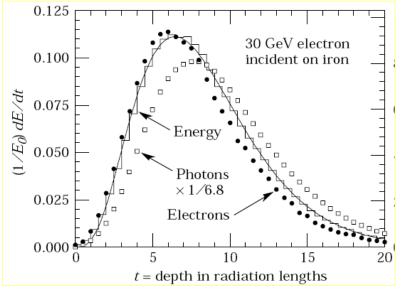
Electromagnetic showers

The shower development is a statistical process.

- α Active sampling wrt total detector volume
- β Uniformity of the detector, non-linearities

Energy resolution

$$\frac{\sigma_E}{E} \approx \frac{\alpha}{\sqrt{E}} \oplus \beta$$

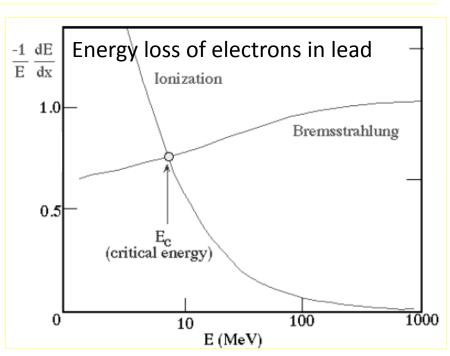


Critical energy

$$\left. \frac{dE}{dx} (E_C) \right|_{Brems} = \frac{dE}{dx} (E_C) \right|_{ion}$$

Shower maximum depth

$$t_{\text{max}} = \frac{\ln\left(\frac{E_0}{E_C}\right)}{\ln 2}$$



Hadronic Showers

Nuclear evaporation

Hadronic interaction:

Elastic:

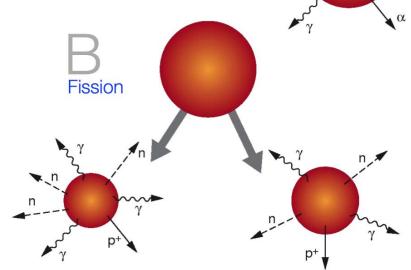
$$p + \text{Nucleus} \rightarrow p + \text{Nucleus}$$

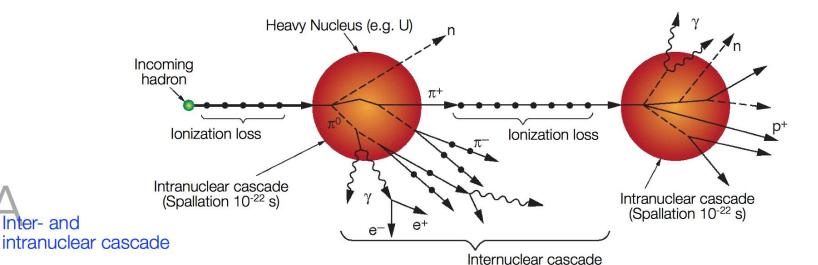
Inelastic:

Inter- and

$$p + \text{Nucleus} \rightarrow \pi^+ + \pi^- + \pi^0 + \dots + \text{Nucleus}^*$$

Nucleus*
$$\rightarrow$$
 Nucleus A + n, p, α , ... \rightarrow Nucleus B + 5p, n, π , ... \rightarrow Nuclear fission



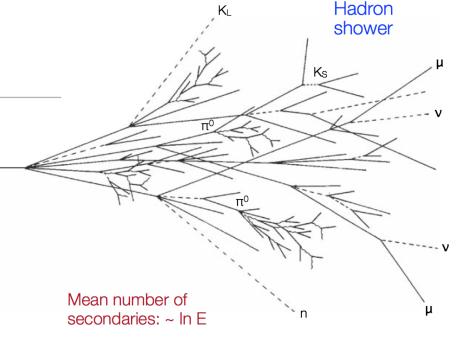


Hadronic Showers

Shower development:

- 1. $p + Nucleus \rightarrow Pions + N^* + ...$
- Secondary particles ...
 undergo further inelastic collisions until they
 fall below pion production threshold
- 3. Sequential decays ...

 $\pi_0 \rightarrow \gamma \gamma$: yields electromagnetic shower Fission fragments $\rightarrow \beta$ -decay, γ -decay Neutron capture \rightarrow fission Spallation ...



Typical transverse momentum: pt ~ 350 MeV/c

Substantial electromagnetic fraction

fem ∼ In E
[variations significant]

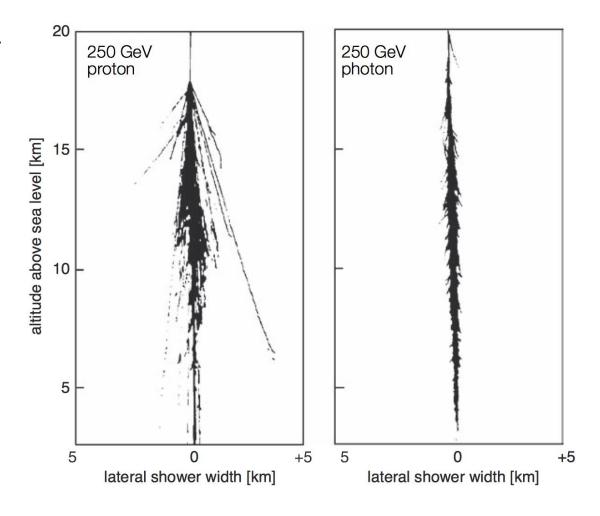
Cascade energy distribution:

[Example: 5 GeV proton in lead-scintillator calorimeter]

5000 MeV [29%]

Hadronic Showers

Comparison hadronic vs electromagnetic shower ... [Simulated air showers]



6. ATOMIC AND NUCLEAR PROPERTIES OF MATERIALS

Table 6.1 Abridged from pdg.lbl.gov/AtomicNuclearProperties by D. E. Groom (2007). See web pages for more detail about entries in this table including chemical formulae, and for several hundred other entries. Quantities in parentheses are for NTP (20°C and 1 atm), and square brackets indicate quantities evaluated at STP. Boiling points are at 1 atm. Refractive indices n are evaluated at the sodium D line blend

Material	Z	A	$\langle Z/A \rangle$		Nucl.inter. length λ_I	Rad.len. X_0	$dE/dx _{\min}$ { MeV	Density $\{ g \text{ cm}^{-3} \}$	Melting point	Boiling point	Refract. index
				_	$\{g \text{ cm}^{-2}\}$	_			(K)	(K)	(@ Na D)
$\overline{\mathrm{H_2}}$	1	1.00794(7)	0.99212	42.8	52.0	63.04	(4.103)	0.071(0.084)	13.81	20.28	1.11[132.]
D_2	1	2.01410177803(8)	0.49650	51.3	71.8	125.97	(2.053)	0.169(0.168)	18.7	23.65	1.11[138.]
He	2	4.002602(2)	0.49967	51.8	71.0	94.32	(1.937)	0.125(0.166)		4.220	1.02[35.0]
Li	3	6.941(2)	0.43221	52.2	71.3	82.78	1.639	0.534	453.6	1615.	
${ m Be}$	4	9.012182(3)	0.44384	55.3	77.8	65.19	1.595	1.848	1560.	2744.	
C diamond	6	12.0107(8)	0.49955	59.2	85.8	42.70	1.725	3.520			2.42
C graphite	6	12.0107(8)	0.49955	59.2	85.8	42.70	1.742	2.210			
N_2	7	14.0067(2)	0.49976	61.1	89.7	37.99	(1.825)	0.807(1.165)	63.15	77.29	1.20[298.]
O_2	8	15.9994(3)	0.50002	61.3	90.2	34.24	(1.801)	1.141(1.332)	54.36	90.20	1.22[271.]
F_2	9	18.9984032(5)	0.47372	65.0	97.4	32.93	(1.676)	1.507(1.580)	53.53	85.03	[195.]
Ne	10	20.1797(6)	0.49555	65.7	99.0	28.93	(1.724)	1.204(0.839)	24.56	27.07	1.09[67.1]
Al	13	26.9815386(8)	0.48181	69.7	107.2	24.01	[1.615]	$2.\dot{6}99$	933.5	2792.	
Si	14	28.0855(3)	0.49848	70.2	108.4	21.82	1.664	2.329	1687.	3538.	3.95
Cl_2	17	35.453(2)	0.47951	73.8	115.7	19.28	(1.630)	1.574(2.980)	171.6	239.1	[773.]
\mathbf{Ar}	18	39.948(1)	0.45059	75.7	119.7	19.55	(1.519)	1.396(1.662)	83.81	87.26	1.23[281.]
${ m Ti}$	22	47.867(1)	0.45961	78.8	126.2	16.16	$1.477^{'}$	$4.\overline{5}40$	1941.	3560.	
${ m Fe}$	26	55.845(2)	0.46557	81.7	132.1	13.84	1.451	7.874	1811.	3134.	
Cu	29	63.546(3)	0.45636	84.2	137.3	12.86	1.403	8.960	1358.	2835.	
Ge	32	72.64(1)	0.44053	86.9	143.0	12.25	1.370	5.323	1211.	3106.	
Sn	50	118.710(7)	0.42119	98.2	166.7	8.82	1.263	7.310	505.1	2875.	
$\mathbf{X}\mathbf{e}$	54	131.293(6)	0.41129	100.8	172.1	8.48	(1.255)	2.953(5.483)	161.4	165.1	1.39[701.]
\mathbf{W}	74	183.84(1)	0.40252	110.4	191.9	6.76	$\stackrel{\cdot}{1.145}^{'}$	19.300	3695.	5828.	
Pt	78	$195.08\dot{4}(9)$	0.39983	112.2	195.7	6.54	1.128	21.450	2042.	4098.	
$\mathbf{A}\mathbf{u}$	79	196.966569(4)	0.40108	112.5	196.3	6.46	1.134	19.320	1337.	3129.	
Pb	82	207.2(1)	0.39575	114.1	199.6	6.37	1.122	11.350	600.6	2022.	
U	92	[238.02891(3)]	0.38651	118.6	209.0	6.00	1.081	18.950	1408.	4404.	
Air (dry, 1 atm) 0.49919			61.3	90.1	36.62	(1.815)	(1.205)		78.80		
Shielding concrete 0.50274			65.1	97.5	26.57	1.711	2.300				
Borosilicate glass (Pyrex) 0.49707			64.6	96.5	28.17	1.696	2.230				

0.42101

0.50000

95.9

66.8

158.0

101.3

7.87

26.54

1.255

1.688

6.220

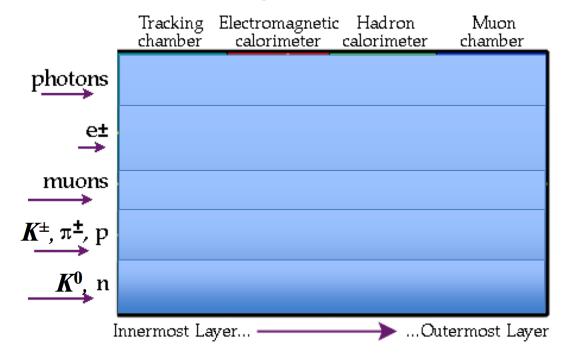
2.650

Lead glass

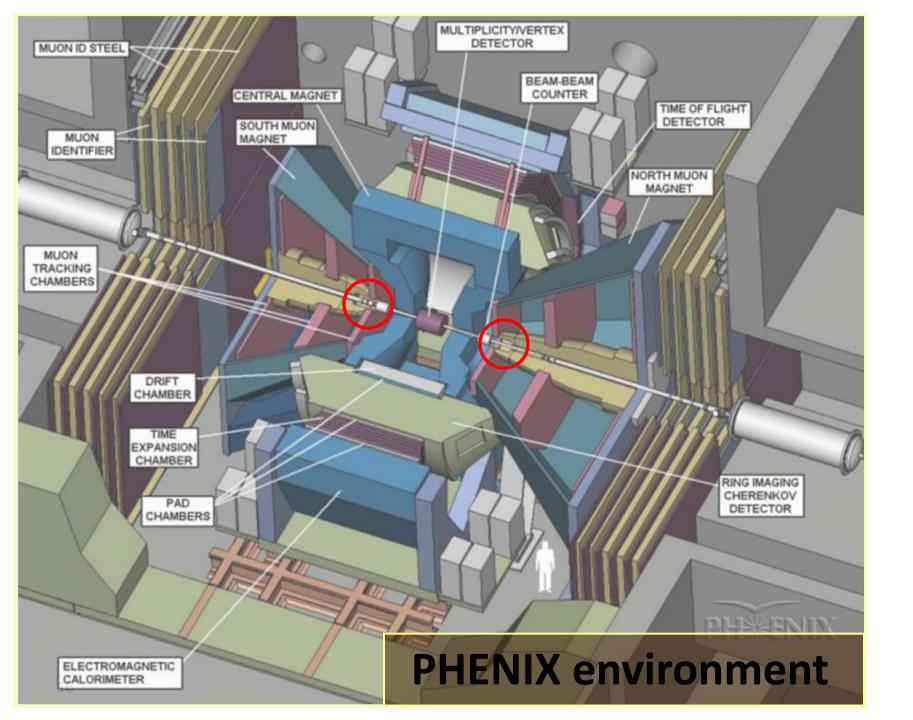
Standard rock

Interactions with Matter

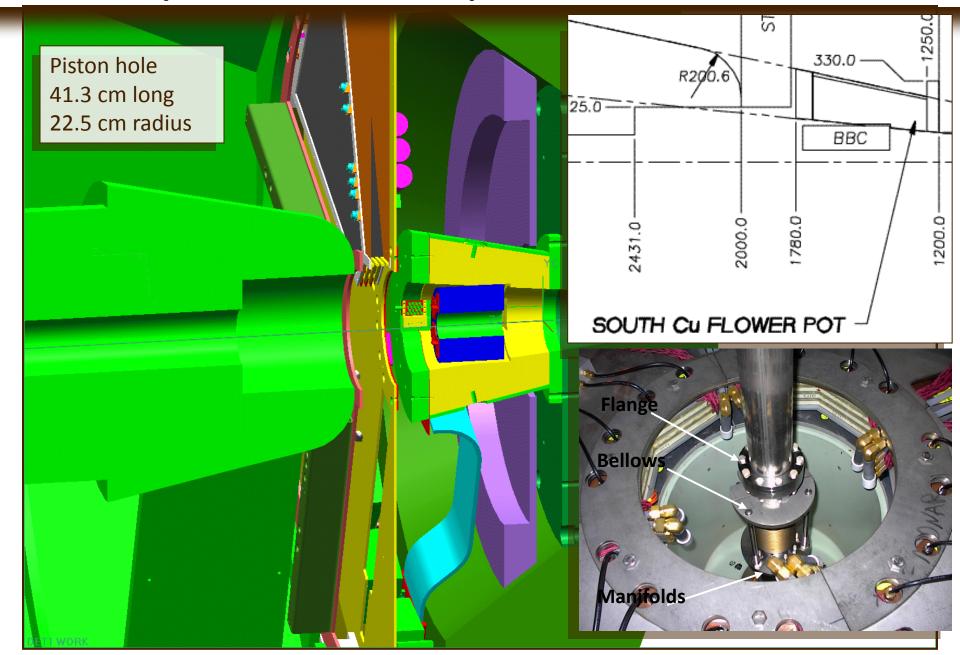
- Generally a detector consists of a tracking detector, an electromagnetic and hadron calorimeter and muon chambers in a magnetic field.
- Each experiment uses different technologies to construct the sub-detectors.



- Tracking detectors measures charged particle trajectory and momentum
- Calorimeter layers measure energy by fully absorbing the particles (destructive measurement).
- Muons do not interact in calorimeter very much: outermost detector to identify muons.

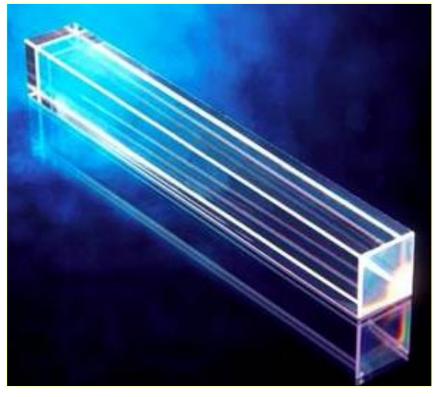


Space in muon piston holes!



PbWO₄ (Lead Tungstate, PWO)

Density	8.28 g/cm ³				
Size	2.2 x 2.2 x 18 cm ³				
Length	20 X ₀ , 0.92 λ				
Weight	721.3 g				
Moliere radius	2.0 cm				
Radiation Length	0.89 cm				
Interaction Length	22.4 cm				
Light Yield	~10 p.e./MeV @ 25° C				
Temp. Coefficient	-2% / °C				
Radiation Hardness	1000 Gy				
Main Emission Lines	420-440, 500 nm				
Refractive Index	2.16				

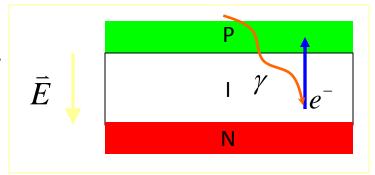


about 50 years in PHENIX forward directions

Avalanche photodiodes?

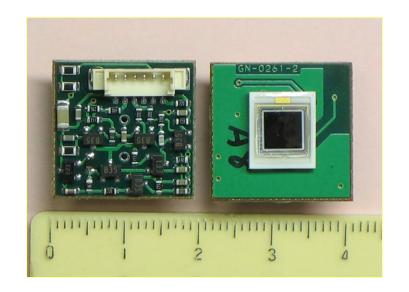
Even small PMTs are sensitive to magnetic fields or expensive (500 – 5000 gauss longitudinally in piston holes)

PIN diodes in reverse bias mode \rightarrow depleted i-layer

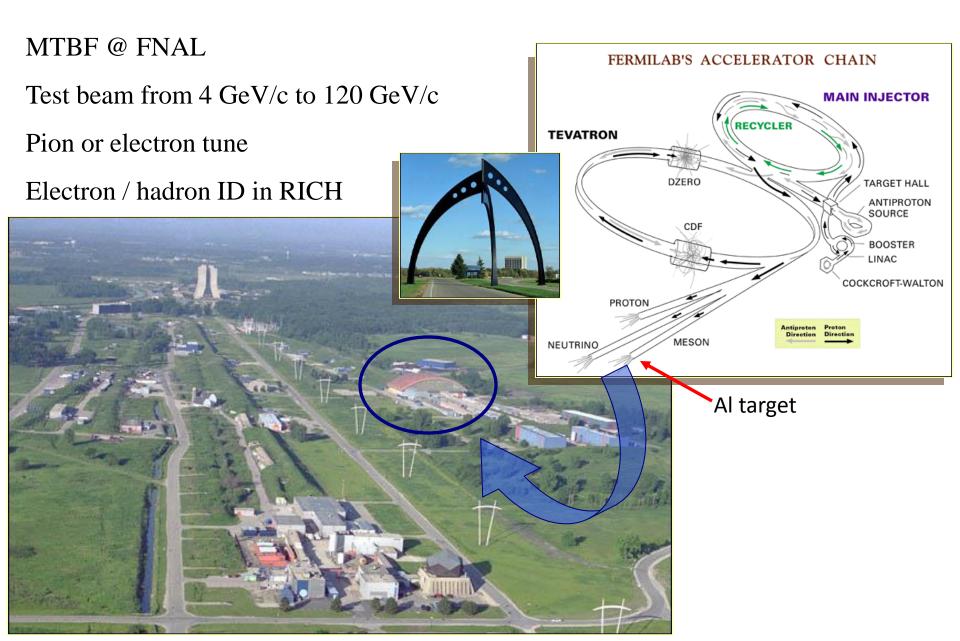


Large reverse bias voltage:
e⁻ acceleration

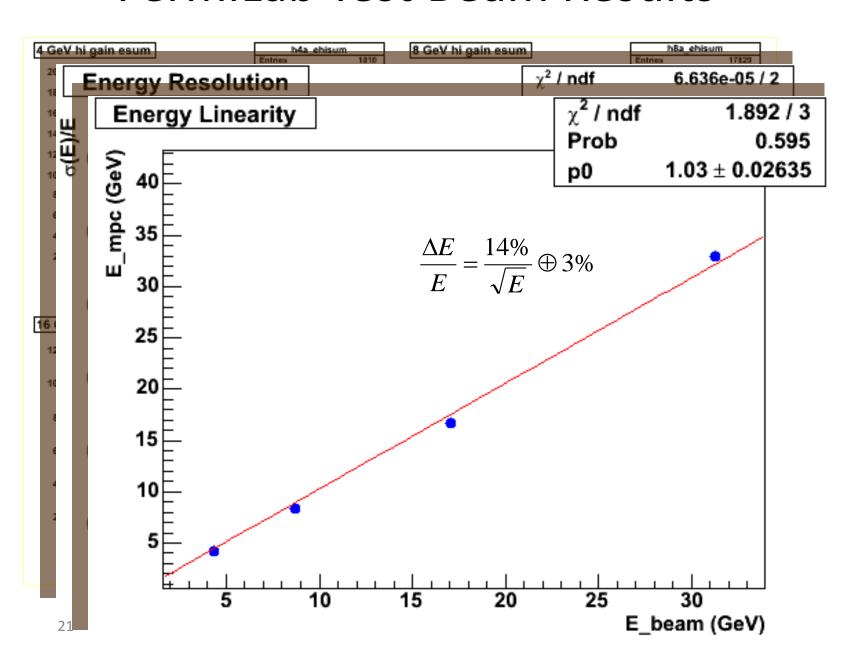
→ collisions with electrons
avalanche multiplication
avalanche leaves the active area



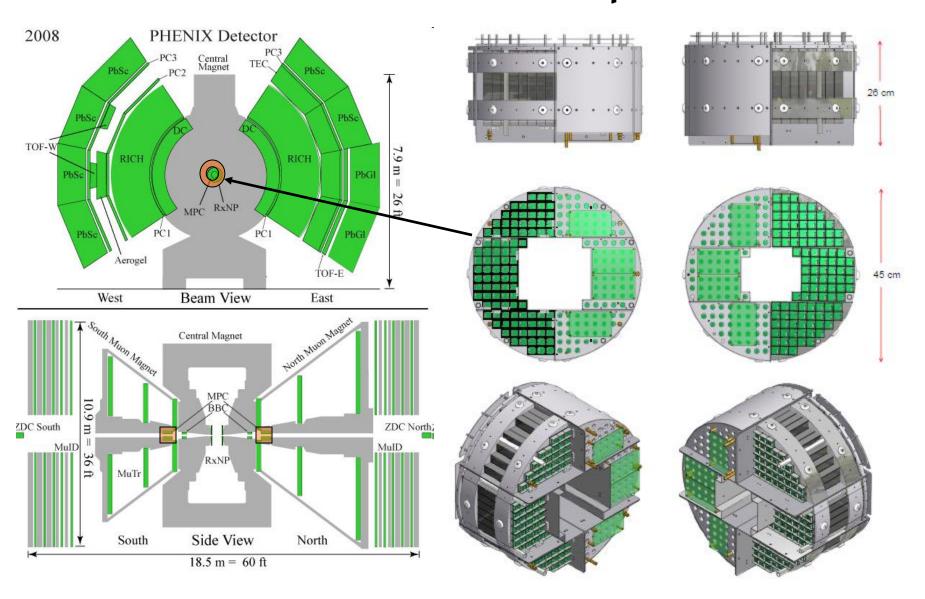
Test beam measurements



FermiLab Test Beam Results



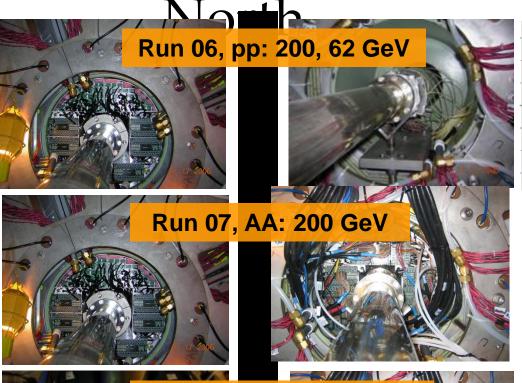
The MPC: $3.1 < |\eta| < 3.8$

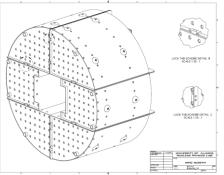


MPC History

South

Initial Installation 192 Towers





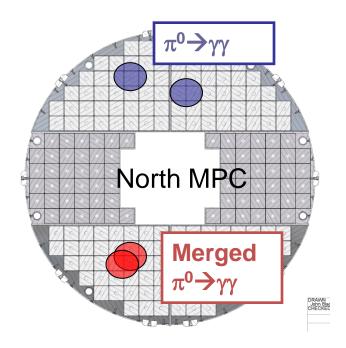
New Detector: 220 Towers

Upgrade: 196 Towers New Monitoring System (N+S)





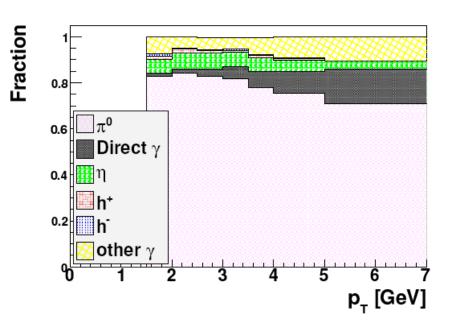
MPC Limitation

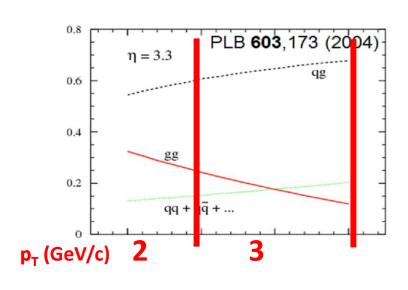


- ID π^0 s up to E ~ 20 GeV with MPCs (3.1 < $|\eta|$ < 3.8)
 - Limitations: tower separation and merging effects
 - Use π⁰s for 7 GeV < E < 22 GeV
 → p_T max ~ 2 GeV/c
- Single Clusters for E > 15 GeV
 - Dominated by π^0 (~ 85%)
 - Access higher p_T

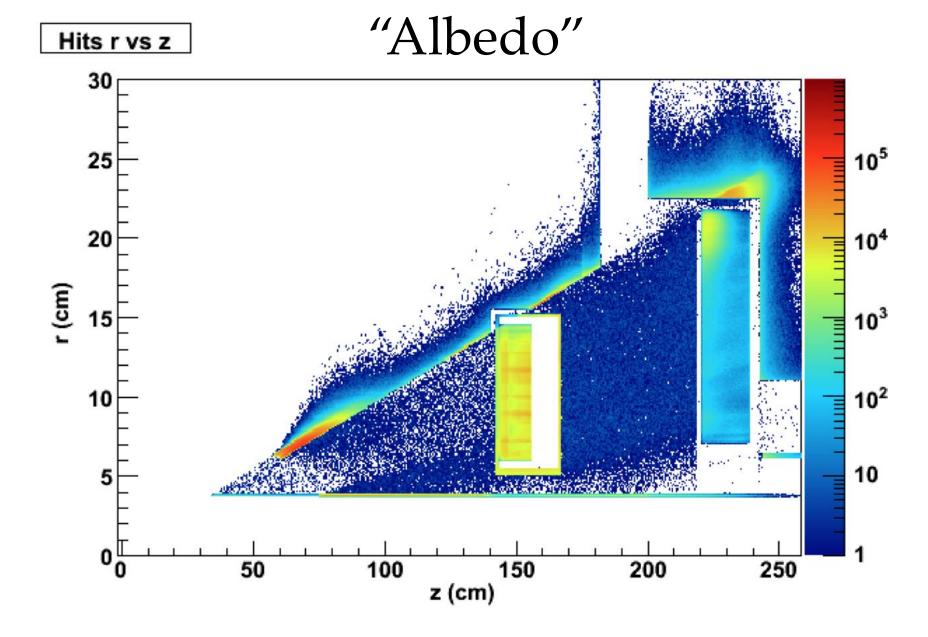
Single Cluster Asymmetry

Cluster Decomposition: Dominated by merged π^{0} 's.





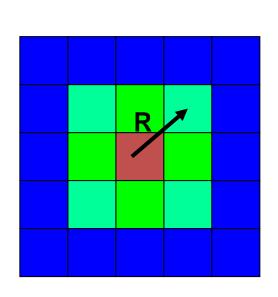
- ✓ Process decomposition skewed more heavily toward quark-gluon than mid-rapidity.
- ✓ Rel lumi uncertainty is significant compared to statistical.

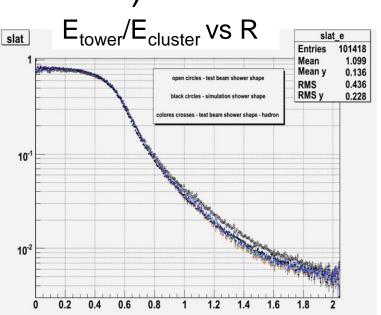


MPC is in a very messy environment! Over a radiation length in front of it.

Reconstruction Algorithm (AN949)

- Basic idea: EMCal clusters = local maxima + surrounding towers
- Use log-weighted positions (weight is 4 + In(E_i/Etot))
- Adapted existing EMCal code to MPC
- Parameterize shower shape and fluctuations from simulation (matches beam test well)

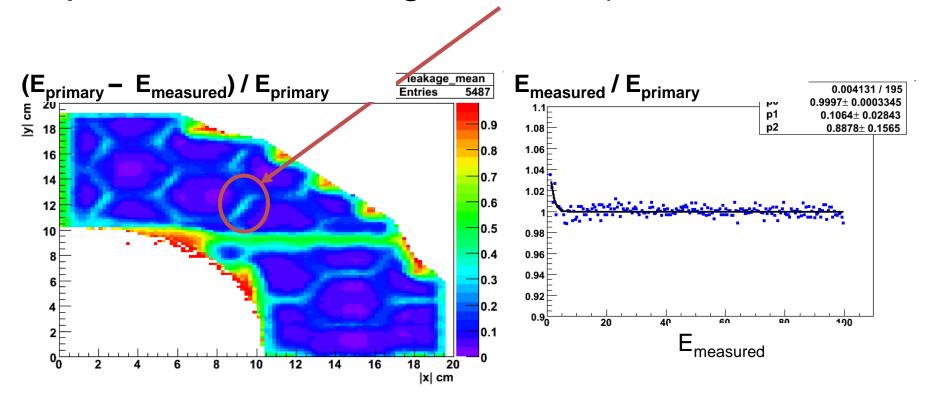




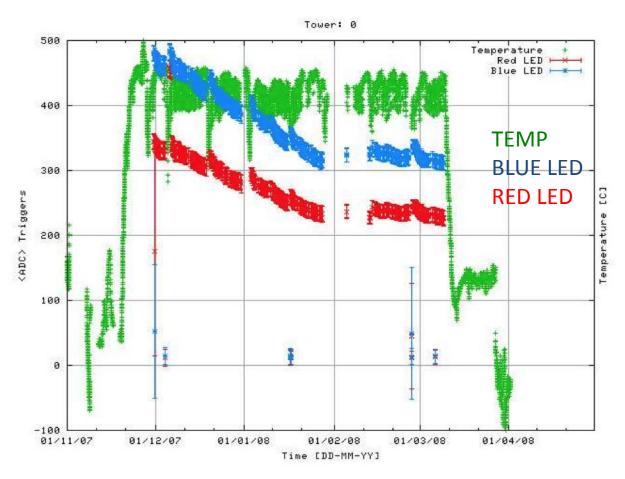
R = distance from tower center to cluster center

MPC Energy Response

 MPCs sit directly behind the BBCs (see shadow below from E = 20 GeV single photons run through GEANT)



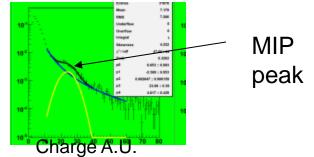
MPC Gain Drop During Run



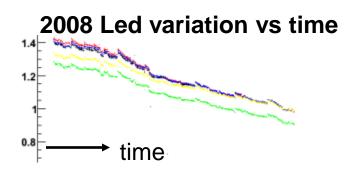
- PbWO4 and APD gains both sensitive to temperature
- PbWO4 suffers massive radiation damage
 - Recovers partially between runs
- APD also suffers neutron damage not recoverable
- LED essential to correcting for this gain drop and fluctuations in gain

Calibrations (our 3rd iteration)

- Use Minimum ionizing particles as first calibration (MIPs deposit 0.234 GeV/tower)
 - Also can use inverse slopes

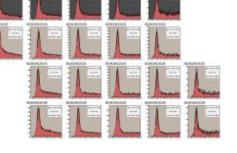


 Correct time-dependence with LEDs (40% over 2008 d+Au, p+p runs)

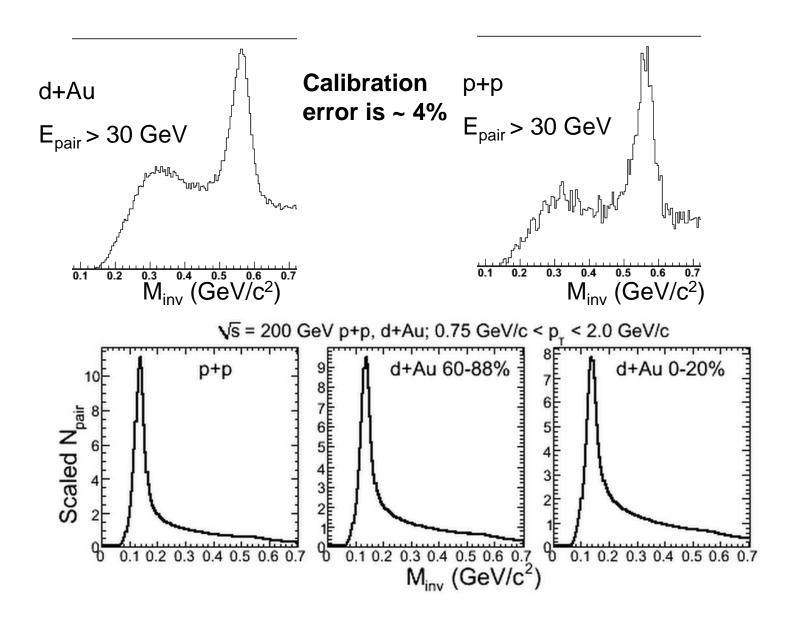


- Use iterative π^0 calibration
 - Match p+p pythia → GEANT simulation masses in each tower

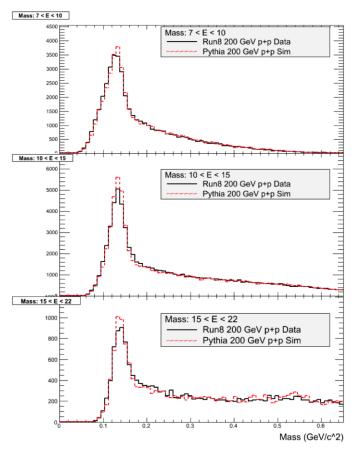
Tower by tower π^0 mass peaks

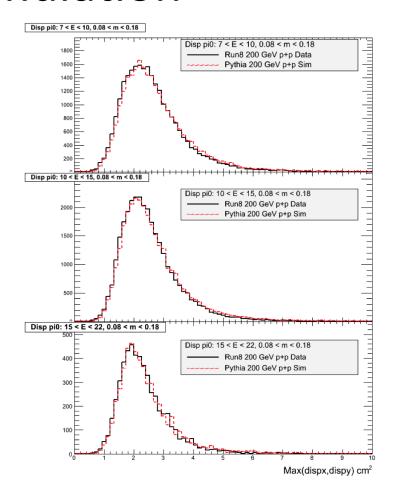


Tests of Calibration: η and π^0 mesons



Tests of Simulation



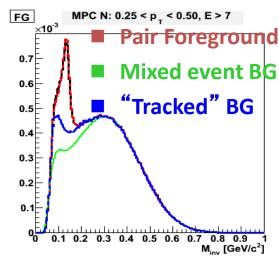


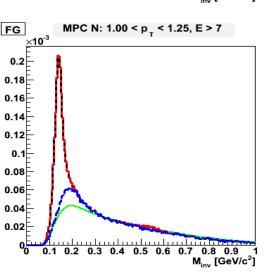
- Figure 7: North MPC invariant mass vs. energy
- Simulation should match the data if one wants to use the simulation for correction factors
- North MPC, Run08 p+p

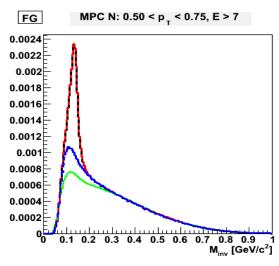
Understanding the Invariant Mass Spectra

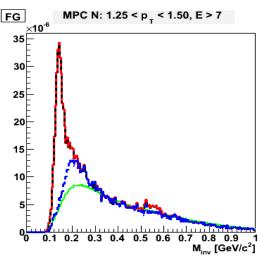
- Developed tools to understand invariant mass spectra
- Track energy depositions for all particles into clusters
- Tracking available in mpcClusterContentV2

 2 embedded p+p Pythia events → PISA →DST

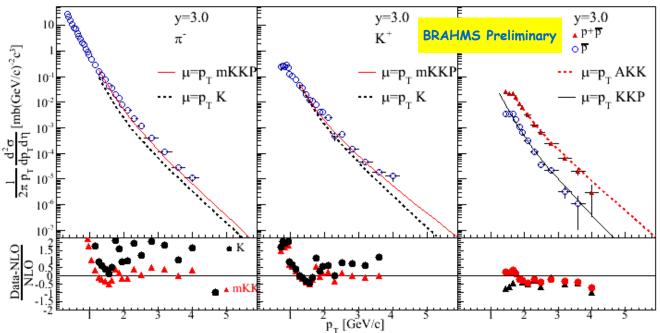




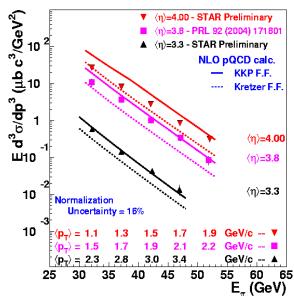




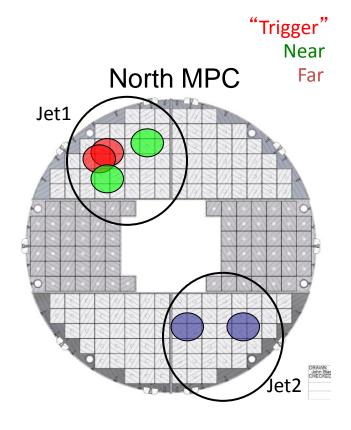
Cross-sections at RHIC, Forward Rapidities

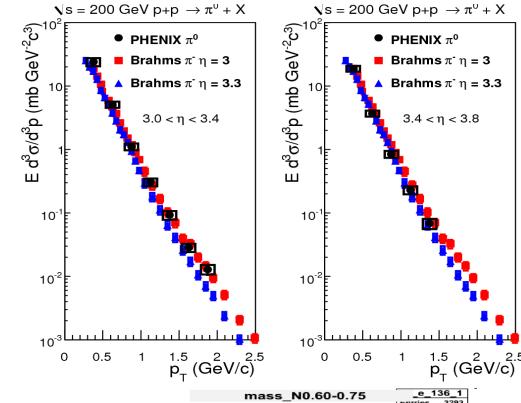


- •Cross-sections generally described well by NLO pQCD at \sqrt{s} = 200 GeV and forward rapidities
- •Are we in a situation where in unpolarized the theory is relatively well understood, but the polarized gives surprises?
 - •Potentially we are in a region where the polarized data gives us new information about QCD, in a region where one can have quantitative theoretical understanding of the effects, and not just phenomenology.



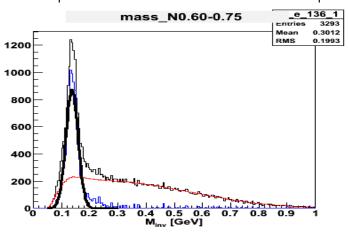
MPC Performance



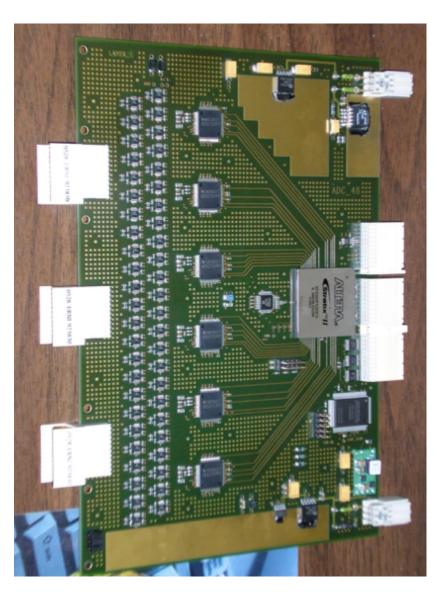


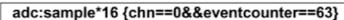
Decay photon impact positions for low and high energy π^0 s. The decay photons from high energy π^0 s merge into a single cluster

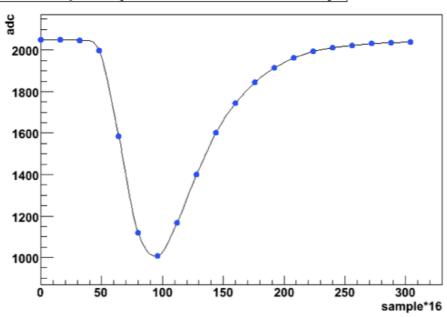
Clusters \geq 80% π^0 (PYTHIA)



Upgraded Electronics in Run12



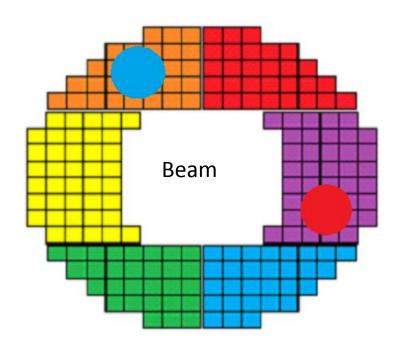


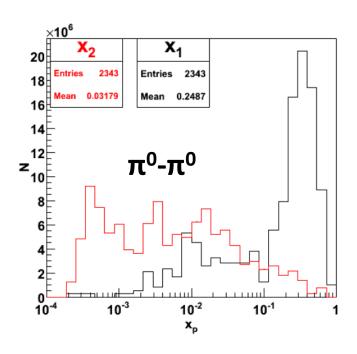


- New electronics in Run12 (HBD ADC boards)
- Digital trigger
 - pT threshold
 - Online gain corrections
 - Remove single tower backgrounds
- Measure energy beyond ADC saturation
- Controls pileup effects

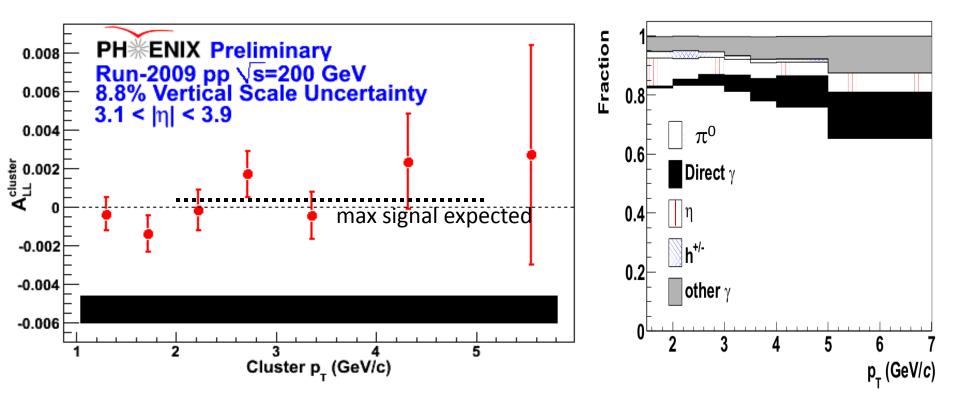
Trigger on Di-hadrons

- MPC now has 6 independent fully digital trigger calculations, arranged azimuthally
- Easy to select for di-hadrons → Increased rejection power
- Allows us to maximize our data purity
- Constrain ΔG at low-x, and with less inclusive probe.



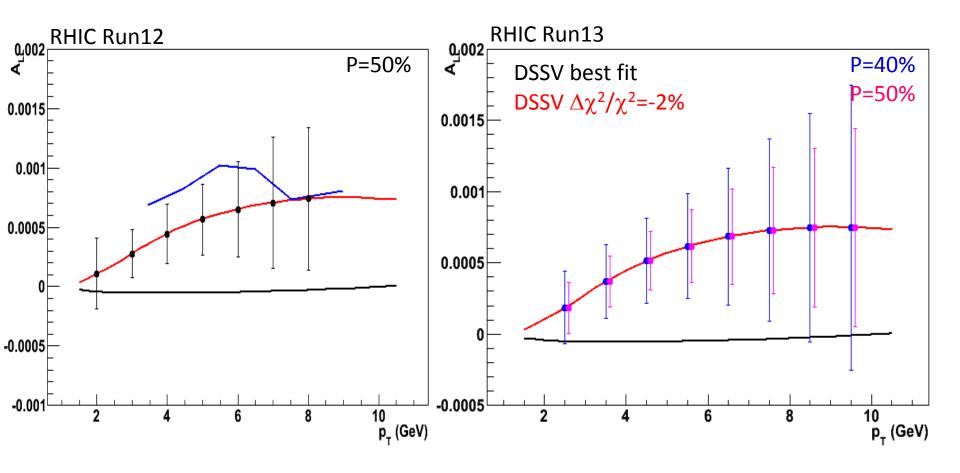


First Forward Measurement of A



- •High p_T EM Cluster Asymmetry, forward pseudo-rapidity 3.1< $|\eta|$ <3.9
 - >80% Merged π^0
- •510 GeV Datasets: Run09, Run11, Run12, Run13
 - •Run12 and Run13 use new MPC electronics with ~4x higher purity

Run12, Run13 Projections



- •Projected error bars around $\delta A_{II} \sim 10^{-4}$
- •From real data, with new MPC Trigger Electronics
- Relative Luminosity Analysis still in progress

Can Clusters be Used?

- 1. Assign a systematic error
 - We have a good idea about the particle type composition
 - Be nice to measure eta's and direct photons separately. Pi-zero's already measured by STAR.
 - We can then calculate what kind of an effect these have on the A_LL given certain assumptions about dG
 - Derive a conservative systematic error based on that
- 2. Try to correct back to pi-zero's
 - Similar to above, but we also apply a correction based on some assumed dG, and assign systematic errors to the uncertainty in that correction.
- 3. Give the theorists our acceptance and efficiency for each particle type
 - With this information they can run this acc*eff filter through their analysis.